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Application of Adomian decomposition method to free vibration analysis of thin isotropic rectangular plates submerged in fluid

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Abstract

Investigating dynamic free vibration response of isotropic thin rectangular plate resting on two-parameter elastic foundation is of interest in the field of geotechnics, structure, highway, railway and mechanical engineering. This study employs analytical approach to the investigation of free vibration analysis of isotropic rectangular plate submerged in fluid and resting on combined elastic foundations under different support conditions. The nonlinear partial differential equation of the system is transformed into nonlinear ordinary differential equation using the Galerkin method of separation. The resulting ordinary differential equation is solved using Adomian decomposition method (ADM). The reliability of the solutions obtained is validated with numerical and existing results as reported in the literature. Also, the bending moment and stress are analysed. The analytical solutions are used for the investigation of dynamic behaviour of plates in fluid, effect of aspect ratio, effect of elastic foundation parameters on natural frequency and effect of bending moment and stress on the mode of the vibrating plate. From the results, it is observed that the increase in elastic foundation parameter increases the natural frequency. Increasing the aspect ratio increases the natural frequency. Increasing the combined elastic foundation parameters increases the natural frequency. Natural frequency of plates reduces when submerged in fluid with the mode shapes not significantly affected. The node and antinodes of the mode shapes are affected by moment and stress. It is expected that the present study will add to the existing knowledge in the field of vibration of plates.

Keywords: Dynamic analysis, Rectangular plate, Deflection, Winkler and Pasternak, Adomian decomposition method

Introduction

Recently, interest has arose into the study of dynamic behaviour of thin rectangular plate. This is because plates are considered an important structural member due to their wide acceptable usage in various fields of engineering. This is also due to their light weight and high strength. Therefore, the study and analyses of natural frequency and modal behaviour are justified. In the study of free vibration analysis of thin rectangular plate on elastic foundation, Motaghian and Akin [1] used an exact method in obtaining the solution. In another study, Rezaiefar and Galal [2] worked on free vibration of rectangular

plate with nonlinear load using finite element method. The adoption of two-parameter foundation has proven to be more reliable than only Winkler foundation due to the ability to take care of the shear interaction among spring element. However, Shariyat and Alipour [3] used a semi-analytical method of solution to investigate the free vibration analysis of functionally graded plate on two-parameter elastic foundation. Meanwhile, Benferhat et al. [4] determine porosity effect on free vibration analysis and bending of functionally graded plate resting on two-parameter foundation. Much attention is given to plates under fluidic interaction due to the importance and application in ship building, marine, nuclear and ocean engineering. The study of plate-fluid interaction in engineering is justified for safety and design purpose. Many literatures had studied the characteristic of immersed and submerged plate in fluids. Lamb [5] carried out an analytical approach into the investigation of fluid-plate coupled system. He determined the natural frequency of clamped circular plate in contact with water using Rayleigh's method; however, Sobamowo et al. [6] investigated plates in a porous medium. In another work, Kwak [7] determined the added virtual mass incremental (AVMI) factor which ascertained the fact that the increase in inertial is due to the presence of fluid in plates. Further study, Soedel and Soedel [8] established a coupled equation of motion of plate carrying fluid. Kerboua et al. [9] later developed a mathematical model for free vibration of plates in contact with fluid using the finite element method. In a later work, Ozdemir [10] employed the finite element method in the vibration response of rectangular Mindlin plate on the Winkler foundation. Also, Mohd et al. [11] performed a numerical analysis on semi-empirical relations. Nikbakhat and Behnamfar [12] demonstrated an experiment on the structures under subway-induced vibrations. In recent time, researches have proposed and applied several semi-analytical methods in analysing dynamic behaviour of rectangular plate resting on nonlinear foundations. Bayat et al. [13] applied the homotopy perturbation method (HPM) for nonlinear free vibration of tapered beams. Werfalli and Aburuga [14] used the Galerkin method for the analysis of rectangular plate. Analytical method [15–18] has proven to be the most reliable method of solution when compared to others like numerical and semi-analytical, but the shortcomings of handling nonlinear problem have not been overcome. Adomian and Rach [19] proposed the method of Adomian decomposition method (ADM) in the beginning of 1980. The aim of ADM polynomial is to weaken the nonlinear terms in the governing equation into a series function. Unlike VIM and HPM that have shortcomings of determination of Lagrange multiplier and small parameters, ADM gives a better solution with fast rate of convergence coupled with high accuracy and minimum calculation because discretization of variables or perturbation is not required: ability to solve nonlinear problems without linearization nor use assumptions of weak nonlinearity. Some related works on ADM are [20–27]. Based on the review of the subject, the topic, analytical approach to investigation of free vibration analysis of isotropic rectangular plate submerged in fluid and resting on combined elastic foundations using ADM has not been investigated. Therefore, the present study focuses on the application of ADM to determine dynamic analysis of isotropic rectangular plates resting on Winkler and Pasternak foundations. The solutions are used for parametric studies.

Problem formulation and mathematical analysis

The present study considers a thin rectangular plate of uniform thickness resting on linear and nonlinear Winkler foundation and Pasternak foundation under different edge

conditions. The boundary edge of the plate may be free, clamped, simply supported or combination of them as shown in Fig. 1. The following assumptions according to Leissa [28] and Chakraverty [29] are made for the development of the governing equation:

1. Normal to the undeformed mid-surface remains straight and normal to the deformed mid-surface with the same length.
2. The thickness of plate is smaller compared to the other dimensions.
3. Rotary inertia and shear deformation effect are negligible.
4. Normal stresses σ_z in the transverse direction to the plate are considered negligibly small.

Similarly, the following assumptions are considered for fluid pressure Δp [9]:

1. Vibration is considered linear; plate is of uniform density and thickness.
2. Flow of fluid is potential, irrotational and homogenous.
3. Fluid is incompressible.
4. Fluid pressure is normal to the plate surface; shear force is ignored because the flow is inviscid.
5. The motion is considered small.
6. The effect of water dynamic loading has an insignificant effect on mode shapes.
7. The system is conservative.

$$\frac{\partial^4 W(x,y)}{\partial X^4} + 2\lambda^2 \frac{\partial^4 W(x,y)}{\partial X^2 \partial Y^2} + \lambda^4 \frac{\partial^4 W(x,y)}{\partial Y^4} - (\Omega^2 + m_{add} - K_L) W(x,y) - K_{NL} W^3(x,y) - g_s \frac{\partial^2 W(x,y)}{\partial X^2} = 0, \tag{1}$$

The governing equation of the thin plate resting linear and nonlinear Winkler and Pasternak foundations also fluid-interaction in dimensionless form is presented in Eq. (1). w is the transverse deflection, where, Ω , g_s and K_L and K_{NL} are the dimensionless natural frequency, dimensionless Pasternak’s Shear stiffness and dimensionless Winkler’s normal stiffness and nonlinear Winkler’s, respectively.

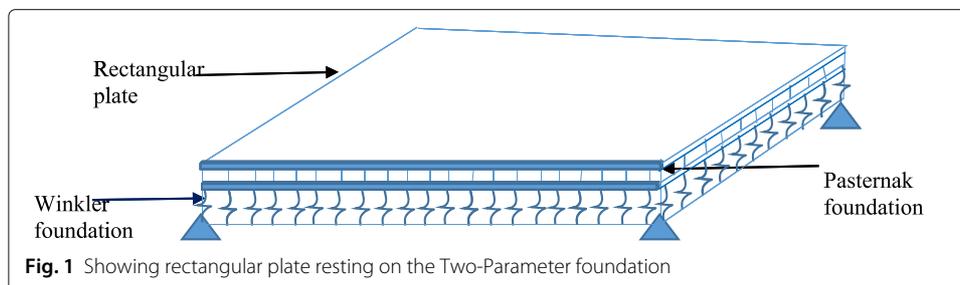
Application of Galerkin method of dimensionless governing equation

Assuming the two opposite edges $Y = 0$ and $Y = 1$ to be simply supported, deflection function can be represented as follows:

$$W = W(X) \sin(m\pi Y), \tag{2}$$

Substitute Eq. (2) into governing differential Eq. (1), we get [22]:

$$\frac{d^4 W(X)}{dX^4} - 2\lambda^2 m^2 \pi^2 \frac{d^2 W(X)}{dX^2} - (\Omega^2 + M_{add} - K_L - \lambda^4 m^4 \pi^4) W(X) - K_{NL} W^3(X) - g_s \frac{d^2 W(X)}{dX^2} = 0, \tag{3}$$



Boundary conditions

Three boundary conditions are considered at $x = 0$ and $x = 1$, notation are used for brevity sake, SS denotes simply supported at both edges of the rectangular plate, SC means simply supported and clamped at the other ans SF denotes simply supported and free edge support. Classical plate theory defines the dimensionless support conditions as follows:

Clamped edge:

$$W = \frac{dW}{dX} = 0, \tag{4}$$

Simply supported edge:

$$W = \frac{d^2W}{dX^2} - \nu (m^2\pi^2\lambda^2) W = 0, \tag{5}$$

Free edge:

$$\frac{d^2W}{dX^2} - \nu (m^2\pi^2\lambda^2) W = 0, \quad \frac{d^3W}{dX^3} - (2 - \nu) (m^2\pi^2\lambda^2) \frac{dW}{dX} = 0, \tag{6}$$

where m is an integer, and ν is the Poisson ratio.

Method of solution: Adomian decomposition method

Description of Adomian decomposition method

The Adomian decomposition method (ADM) is developed by George Adomian in 1990s. It is a semi-analytical method of solving partial and ordinary nonlinear equation. It uses ‘Adomian polynomials’ for fast convergence of nonlinear aspect of differential equation. The principle of operation is as follows:

Considering the following equation:

$$Lw + Nw + Rw = g(x) \tag{7}$$

where L is a linear operator, N is the nonlinear operator, R is the remaining linear operator and g is the inhomogeneous term. If L is the fourth order operator, it is defined by:

$$L = \frac{d^4}{dx^4}, \tag{8}$$

Assuming L is invertible, then inverse operator L^{-1} is given as:

$$L^{-1}(\cdot) = \int_0^x \int_0^x \int_0^x \int_0^x (\cdot) dx dx dx dx, \tag{9}$$

Therefore,

$$L^{-1}Lw = w(x) - w(0) - xw'(0) - \frac{1}{2!}x^2w''(0) - \frac{1}{3!}x^3w'''(0), \tag{10}$$

Applying L^{-1} to both sides Eq. (7) gives:

$$w = \Phi_0 - L^{-1}Rw - L^{-1}Nw + L^{-1}g(x), \tag{11}$$

where

$$\Phi_0 = \begin{cases} w(0), & \text{if } L = \frac{d}{dx} \\ w(0) + xw'(0), & \text{if } L = \frac{d^2}{dx^2} \\ w(0) + xw'(0) + \frac{x^2}{2!}w''(0), & \text{if } L = \frac{d^3}{dx^3} \\ w(0) + xw'(0) + \frac{x^2}{2!}w''(0) + \frac{x^3}{3!}w'''(0), & \text{if } L = \frac{d^4}{dx^4} \end{cases} \tag{12}$$

The decomposition principle comprises decomposing the solution to the sum of infinite number of terms defined by the series:

$$w = \sum_{n=0}^{\infty} w_n, \quad (13)$$

The nonlinear term is written as:

$$N(w) = \sum_{n=0}^{\infty} A_n, \quad (14)$$

where A_n 's are the Adomian polynomials.

$$A_n = \frac{1}{n!} \left[\frac{\partial^n}{\partial \xi^n} \left(\left(\sum_{i=0}^{\infty} \xi^i w_i(x) \right) \right) \right]_{\xi=0} \quad n = 0, 1, 2, 3 \dots; \quad (15)$$

where ξ is a grouping parameter.

$$\begin{aligned} A_0 &= N(w_0), \\ A_1 &= N(w_0)w_1, \\ A_2 &= N'(w_0)w_2 + \frac{1}{2!}N''(w_0)w_1^2, \end{aligned} \quad (16)$$

Substituting Eqs. (13) and (14) into (7) gives:

$$\sum_{n=0}^{\infty} w_n = \Phi_0 - L^{-1}R \left(\sum_{n=0}^{\infty} w_n \right) - L^{-1} \left(\sum_{n=0}^{\infty} A_n \right) + L^{-1}g(x), \quad (17)$$

The iterative schemes are:

$$w = \Phi_0 + L^{-1}g(x), \quad (18)$$

$$w_{n+1} = -L^{-1}Rw_n - L^{-1}A_n, \quad (19)$$

where $n \geq 0$

This results to:

$$\begin{aligned} w_0 &= \Phi_0 + L^{-1}g(x), \\ w_1 &= -L^{-1}Rw_0 - L^{-1}A_0, \\ w_2 &= -L^{-1}Rw_1 - L^{-1}A_1, \end{aligned} \quad (20)$$

where $n \geq 0$

Application of ADM to the solution of nonlinear equation under investigation

According to basic rule of ADM, Eq. (3) is written as:

$$w_n(x) = w(0) + w'(0)x + w''(0)\frac{x^2}{2!} + w'''(0)\frac{x^3}{3!} - L^{-1}[Rw] - L^{-1}[Nw] + L^{-1}[g(x)], \quad (21)$$

Applying Eq. (5) into Eq. (21), we have:

$$\begin{aligned} w_n(x) &= w'(0)x + w'''(0)\frac{x^3}{3!} \\ &- L^{-1} \left[2m^2\pi^2\lambda^2\frac{d^2w}{dx^2} - (\Omega^2 + M_{\text{add}} - k_L - \lambda^4m^4\pi^4)w - g_s\frac{d^2w}{dx^2} \right] - L^{-1}[k_{\text{NL}}w^3], \end{aligned} \quad (22)$$

$w'(0)$ and $w'''(0)$ are unknown which are represented by B and C , respectively:

$$w_n(x) = Bx + \frac{Cx^3}{3!} - L^{-1} \left[2m^2\pi^2\lambda^2 \frac{d^2w}{dx^2} - (\Omega^2 + M_{\text{add}} - k_L - \lambda^4 m^4 \pi^4) w - g_s \frac{d^2w}{dx^2} \right] - L^{-1} [k_{\text{NL}} w^3], \quad (23)$$

In Adomian's polynomial, A 's are expressed as:

$$\begin{aligned} A_0 &= (w_0)^3; \\ A_1 &= 3(w_0)^2 w_1; \\ A_2 &= 3(w_0)^2 w_2 + 3w_0 (w_1)^2; \\ A_3 &= 6w_0 w_1 w_2 + (w_1)^3 + 3(w_0)^2 w_3; \\ A_4 &= 6w_0 w_1 w_3 + 3w_0 (w_2)^2 + 3(w_1)^2 w_2 + 3(w_0)^2 w_4; \end{aligned} \quad (24)$$

The other polynomials are generated in similar way as Eq. (24).

$$w_0 = Bx + \frac{1}{6}Cx^3; \quad (25)$$

$$w_1 = -L^{-1} \left[2m^2\pi^2\lambda^2 \frac{d^2w_0}{dx^2} - (\Omega^2 + M_{\text{add}} - k_L - \lambda^4 m^4 \pi^4) w_0 - g_s \frac{d^2w_0}{dx^2} \right] - L^{-1} [k_{\text{NL}} A_0], \quad (26)$$

$$w_2 = -L^{-1} \left[2m^2\pi^2\lambda^2 \frac{d^2w_1}{dx^2} - (\Omega^2 + M_{\text{add}} - k_L - \lambda^4 m^4 \pi^4) w_1 - g_s \frac{d^2w_1}{dx^2} \right] - L^{-1} [k_{\text{NL}} A_1], \quad (27)$$

$$w_3 = -L^{-1} \left[2m^2\pi^2\lambda^2 \frac{d^2w_2}{dx^2} - (\Omega^2 + M_{\text{add}} - k_L - \lambda^4 m^4 \pi^4) w_2 - g_s \frac{d^2w_2}{dx^2} \right] - L^{-1} [k_{\text{NL}} A_2], \quad (28)$$

\vdots

Using the definition (13), the analytical series solution is obtained.

$$W(x) = \sum_{m=0}^{\infty} W_m(x) = W_0(x) + W_1(x) + W_2(x) + \dots; \quad (29)$$

where the constant B and C are found using Eqs. (4), (5) and (6). Setting the controlling parameters to 0 then substitute the values into Eq. (29) lead to the following simultaneous expression:

$$\begin{aligned} \psi_{11}^{(n)}(\Omega) w_0 + \psi_{12}^{(n)}(\Omega) w_2 &= 0 \\ \psi_{21}^{(n)}(\Omega) w_0 + \psi_{22}^{(n)}(\Omega) w_2 &= 0 \end{aligned} \quad (30)$$

The polynomials ψ_{11} , ψ_{12} , ψ_{21} and ψ_{22} are represented in terms of the natural frequency Ω . Meanwhile, ψ_{11} , ψ_{12} , ψ_{21} and ψ_{22} are representing a series expression obtained from Eq. (29). Therefore, Eq. (30) may be written in matrix form as:

$$\begin{bmatrix} \psi_{11}^{(n)}(\Omega) & \psi_{12}^{(n)}(\Omega) \\ \psi_{21}^{(n)}(\Omega) & \psi_{22}^{(n)}(\Omega) \end{bmatrix} \begin{Bmatrix} w_0 \\ w_2 \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}, \quad (31)$$

The following characteristic determinant is obtained applying the non-trivial condition:

$$\begin{bmatrix} \psi_{11}^{(n)}(\Omega) & \psi_{12}^{(n)}(\Omega) \\ \psi_{21}^{(n)}(\Omega) & \psi_{22}^{(n)}(\Omega) \end{bmatrix} = 0, \quad (32)$$

Solving Eq. (32) gives the natural frequencies.

$$\frac{|\Omega_j^{(i)} - \Omega_j^{(i-1)}|}{|\Omega_j^{(i)}|} \leq \varepsilon, j = 1, 2, 3, \dots, n \tag{33}$$

where the iteration counter is represented by i , the estimated value of the j th dimensionless natural frequency is $\Omega_j^{(i)}$ and small number chosen is ε . For this study, $\varepsilon = 0.0001$. The results show that the use of only few iterations gives a convergent solution.

The deflection series solution of the governing equation for SF boundary condition is:

$$\begin{aligned} W(x) = & x - \frac{8367x^3}{14521} + \frac{24211x^5}{75362} - \frac{8315x^7}{314432} + \frac{1515x^9}{369907} - \frac{23x^{11}}{178680} + \frac{107x^{13}}{11628921} - \frac{5x^{15}}{33008094} \\ & + \frac{4x^{17}}{644113311} - \frac{15929564220}{x^{19}} + \frac{599979099008}{x^{21}} - \frac{87813673294953}{x^{23}} + \frac{4724940354278380}{x^{25}} \\ & - \frac{959420400780756366}{x^{27}} + \frac{69863109901089499689}{x^{29}} - \frac{18793443763260581094762}{x^{31}} \\ & + \frac{1779726318335411523907948}{x^{33}} - \frac{612598163066578778270048320}{x^{35}} + \dots \end{aligned} \tag{34}$$

Modified of Adomian decomposition method procedure

The small domain limitation of semi-analytical method has been overcome by the introduction of an after treatment method in power series method. Laplace-Padè approximant has proven to be a very reliable approach that also increases the convergence rate of the iteration. The Padè is a form of converting the analytical solution obtained through ADM method to polynomial rational form. The basic procedures are as follows:

1. Apply Laplace transform to the series solution Eq. (34) and setting:

$$s = \frac{1}{t}. \tag{35}$$

2. Apply Padè approximation to the solution from previous step to obtain the following approximation in the following rational form:

$$\left[\frac{L}{M} \right] = \frac{P_0 + P_1t + P_2t^2 + \dots + P_Lt^L}{q_0 + q_1t + q_2t^2 + \dots + q_Mt^M}. \tag{36}$$

and setting $t = \frac{1}{s}$.

3. Apply inverse Laplace transform on:

$$\left[\frac{L}{M} \right] \tag{37}$$

approximant.

Application of Adomian-Padè method

The accuracy of the ADM is improved using the principle of Laplace-Adomian-Padè method (MADM)

Apply Laplace transform to the series solution Eq. (34) as:

$$\begin{aligned} \mathcal{L}[w(t)] = & \frac{1}{s^2} - \frac{50202}{14521s^4} + \frac{173212}{4493s^6} - \frac{1268962}{9521s^8} + \frac{1141417}{768s^{10}} - \frac{1253711}{244s^{12}} + \frac{15240749}{266s^{14}} - \frac{13271621}{67s^{16}} \\ & + \frac{55221251}{25s^{18}} - \frac{45818617}{6s^{20}} + \frac{255463610}{3s^{22}} - \frac{294396257}{s^{24}} + \frac{3282837217}{s^{26}} - \frac{11349424550}{s^{28}} \\ & + \frac{126558379800}{s^{30}} - \frac{437537619900}{s^{32}} + \frac{4879018493000}{s^{34}} - \frac{16867742330000}{s^{36}} + \frac{188093601500000}{s^{38}} \\ & - \frac{650277183600000}{s^{40}}, \end{aligned} \tag{38}$$

Also, replacing $s = \frac{1}{t}$, and calculating Padè approximant of $[\frac{5}{5}]$ and letting $t = \frac{1}{s}$, gives the following:

$$\left[\frac{5}{5}\right] = \frac{\left(s^2 - \frac{50202}{14521}\right)}{\left(s^4 - \frac{173212}{4493}\right)}, \tag{39}$$

Applying the inverse Laplace transform to the Padè approximant of Eq. (39), the MADM solution is:

$$w(t) = \frac{76381}{244507} \sin\left(\frac{50800t}{20387}\right) + \frac{12871}{144730} \sinh\left(\frac{50800t}{20387}\right), \tag{40}$$

Moment and stress

The bending moment of M_x and M_y in x and y direction are:

$$M_x = -D_0 \left(\frac{\partial^2 w}{\partial x^2} + \nu \frac{\partial^2 w}{\partial y^2} \right), \tag{41}$$

The surface stresses σ_x and σ_y in the x and y directions are:

$$\sigma_x = 6 \frac{M_x}{h^2}, \tag{42}$$

where D_0 is the flexural rigidity, ν is the Poisson’s ratio and w is the transverse deflection.

Comparison of the results using numerical method

Obtaining analytical solution for the natural frequency of this governing differential non-linear Eq. (1) may be very difficult due to the nonlinear Winkler parameters introduced. Therefore, the next alternative approach is the numerical method. In this study, the finite difference method is employed for the discretization of the governing Eq. (3) along with the boundary conditions of Eqs.(4), (5) and (6). The rectangular plate illustrated in Fig. 2 is meshed along the space x divided into equally spaced (Δx) points. The mesh points in the space are in $x_1, x_2, x_3, \dots, x_{N-1}, x_N$ where the edge nodes are x_1 and x_N , respectively.

$$\Delta x = \frac{x_N - x_0}{N} \Rightarrow \frac{1 - 0}{11} = 0.09$$

For the validation of the analytical solutions obtained from the ADM, controlling parameters are set as 0 for validation. Investigating the effect of nonlinear Winkler parameter, nonlinear equation is formed and solved using the Jacobian matrix of Newton-Raphson method.

Discretizing the governing equation

$$\frac{w_{i+2} - 4w_{i+1} + 6w_i - 4w_{i-1} + w_{i-2}}{(\Delta x)^4} - 2\lambda^2 m^2 \pi^2 \left\{ \frac{w_{i+1} - 2w_i + w_{i-1}}{(\Delta x)^2} \right\} - (\Omega^2 + M_{add} - k_L - \lambda^4 m^4 \pi^4) w_i - k_{NL} w_i^3 - g_s \left\{ \frac{w_{i+1} - 2w_i + w_{i-1}}{(\Delta x)^2} \right\} = 0, \tag{43}$$

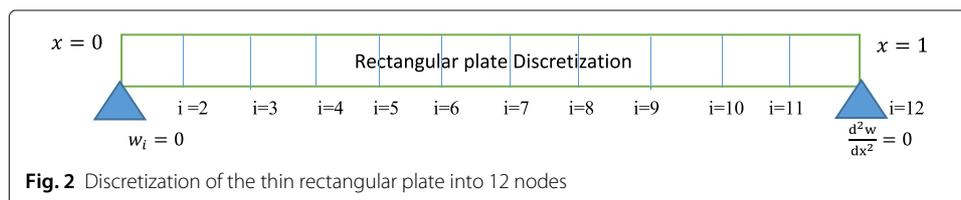


Fig. 2 Discretization of the thin rectangular plate into 12 nodes

Having applied the finite difference approximation, grouping and rearranging the term, the final algebraic form of the finite difference equation becomes:

$$\begin{aligned} & w_{i+2} - 4w_{i+1} + 6w_i - 4w_{i-1} + w_{i-2} - 2\lambda^2 m^2 \pi^2 (\Delta x)^2 (w_{i+1} - 2w_i + w_{i-1}) \\ & - (\Omega^2 + M_{add} - k_L - \lambda^4 m^4 \pi^4) (\Delta x)^4 w_i + k_{NL} (\Delta x)^4 w_i^3 - g_s (\Delta x)^2 \\ & (w_{i+1} - 2w_i + w_{i-1}) = 0 \end{aligned} \quad (44)$$

The discretization in Eq. (43) clearly shows that central difference approximation is employed in the discretization of the governing equation. Indirectly, the inner nodes ($i = 2 : N - 1$) will be utilized for solving the model. Therefore, only ten interior nodes will be used while the two node boundaries are excluded.

Discretizing the boundary conditions

For the sake of brevity, only the SS boundary condition is shown here, while the other boundary conditions SC and SF are solved following the same approach. The starting point and end are presented in the following unique form:

$$x = i, w_i = 0, \frac{d^2 W}{dX^2} = 0 \Rightarrow \frac{w_{i+1} - 2w_i + w_{i-1} - \nu (m^2 \pi^2 \lambda^2) w_i}{(\Delta x)^2} = 0, \quad (45)$$

$i = 1, w_1 = 0, w_2 - 2w_1 + w_0 - \nu (m^2 \pi^2 \lambda^2) w_1 = 0$. Therefore, $w_2 = -w_0$

$$x = i, w_i = 0, \frac{d^2 W}{dX^2} - \nu (m^2 \pi^2 \lambda^2) W = 0 \Rightarrow \frac{w_{i+1} - 2w_i + w_{i-1} - \nu (m^2 \pi^2 \lambda^2) w_i}{(\Delta x)^2} = 0, \quad (46)$$

$i = 12, w_{12} = 0, w_{13} - 2w_{12} + w_{11} - \nu (m^2 \pi^2 \lambda^2) w_{12} = 0, w_{13} + w_{11} = 0$. Therefore,

$$i \Rightarrow 2 \quad w_4 - \frac{28125w_3}{6761} + \frac{42728w_2}{6761} + w_0 - \frac{(-639\pi^4 + 639\Omega^2) w_2}{9739369} = 0, \quad (47)$$

\vdots

$$i \Rightarrow 11 \quad w_{13} + \frac{42728w_{11}}{6761} - \frac{28125w_{10}}{6761} + w_9 - \frac{(-639\pi^4 + 639\Omega^2) w_{11}}{9739369} = 0, \quad (48)$$

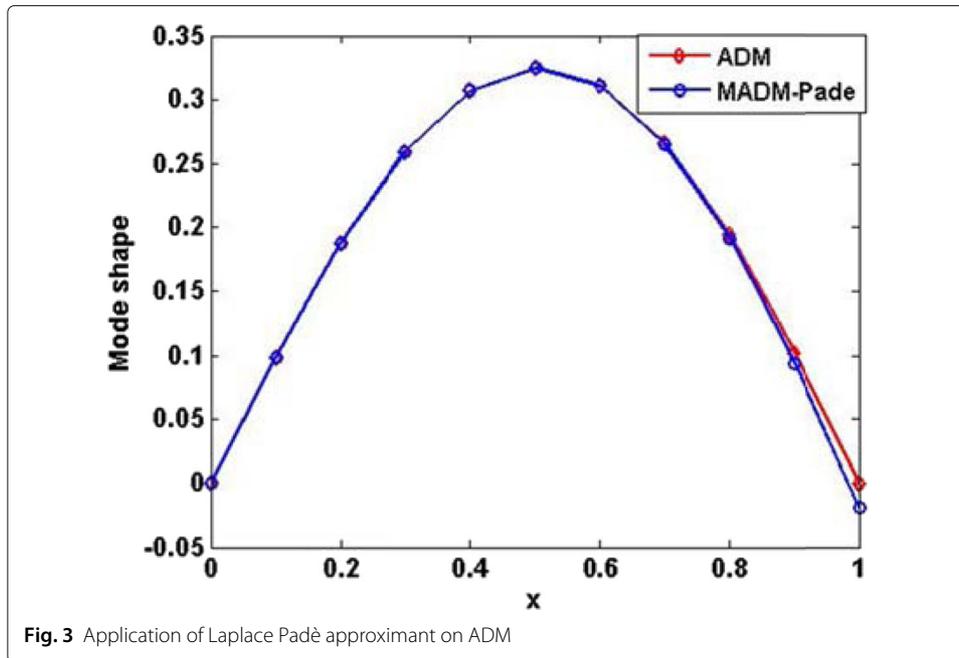
Resolving the matrix yields the natural frequency.

Results and discussion

ADM is used in obtaining solution for free vibration analysis of a uniform thickness thin rectangular plate resting on elastic foundation and submerged in fluid. The accuracy of the analytical solutions obtained is validated with numerical and analytical solution of different method as reported in the cited literature [30] and presented in Table 1. Good agreement of the result is observed along the entire values under different boundary conditions. Generally, the natural frequency is expressed in dimensionless form Ω . Comparing the analytical solution with the numerical shows that even for 5 iterations compared to 10 of numerical, the hybrid method still predicts better than the numerical, though the numerical has an edge handling the nonlinear term in the model. Results obtained with ADM is equally in good agreement compared to the analytical solution obtained by the reported literature [31] using DTM. The analytical solution by DTM converges at 21 iterations compared to ADM of 5 iterations. It is observed that the number of iterations needed to obtain convergence in relation to the natural frequency varies. For instance, fundamental mode requires 2–3 iterations for ADM while the second and third

Table 1 Illustrate the validation of the analytical solution

Edge condition/natural frequency	Simply supported (SS)	Present FDM	Present ADM	Simply supported clamped (SC)	Present FDM	Present ADM	Simply supported free (SF)	Present FDM	Present ADM
Ω_1	Bhat [32] 19.7392	19.8713	19.7434	Bhat [32] 23.6463	23.6358	23.3324	Kumar [31] 11.7195	11.7865	11.7445

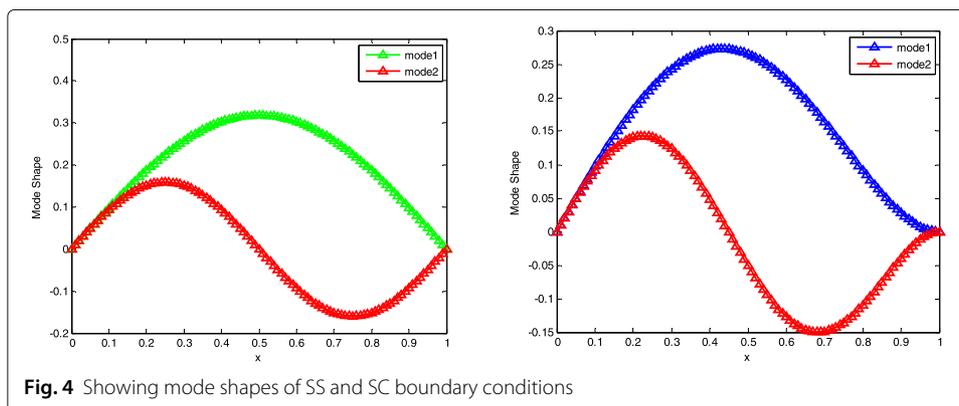


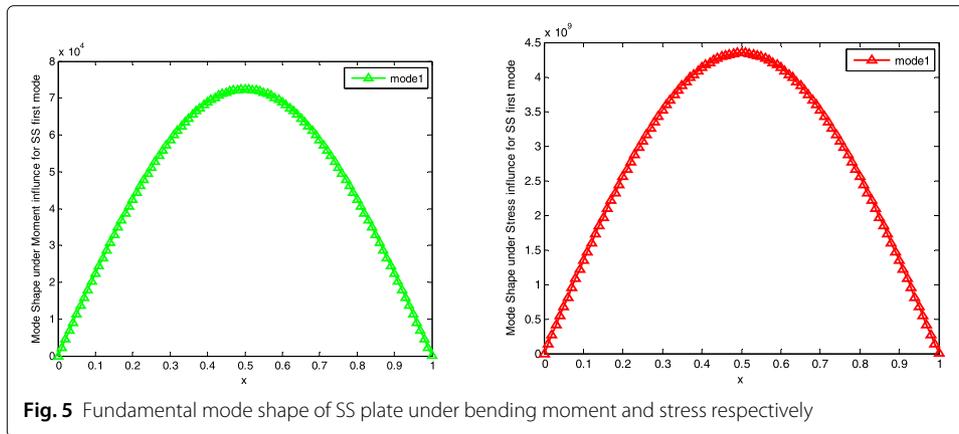
iterations require more iterations. This phenomenon is attribute to more complex series function combination. The results shown in Table 1 demonstrate that, though fundamental natural frequency gives a reasonable prediction of the rectangular plate behaviour, more iterations are still required to give other higher mode natural frequencies and also increase the accuracy. It is interesting to note that present results agree very well with the past results.

Figure 3 illustrates the modification of the Adomian decomposition method by using the Laplace-Padè approximant approach. It is observed that MADM helps to improve the results most especially when the domain of the analysis is getting wider.

Mode shapes

The mode shape for the first two natural frequencies are shown in Fig. 4. It is important to note that the mode shape obeys the classical theory of vibration. For bending moment and stress, location of the vibrating node and antinodes are somehow different due to the vanishing mode of the boundary condition. Figures 4 and 5 show mode shape due to the bending moment and stress; it is clearly shown that the location of node and antinodes of





the vibrating plate changes. While Fig. 6 depicts the extrema mode shape variation. It is clearly revealed that Figs. 5 and 6 when compared to Fig. 4 have different mode shapes. Invariably, the extrema mode shape location differs based on the boundary conditions. When the plate is submerged in fluid, the same mode shapes as illustrated in Fig. 4 are obtained; this is in accordance with the experimental result obtained by [33, 34]

Variation of foundation parameter on natural frequency

To investigate the variation effect of the elastic foundation on the first two natural frequencies of the thin rectangular plate under different boundary conditions discussed earlier, the material properties of the thin rectangular plate model are as follows [9]:

$E = 207 \text{ Gpa}$, material density $\rho = 7850 \text{ kg/m}^3$, thickness of the plate $h = 0.01 \text{ m}$ and density of the water $\rho = 1000 \text{ kg/m}^3$. Thin rectangular plate is considered. Reservoir with a dimension of $5 \text{ m} \times 5 \text{ m} \times 5 \text{ m}$ is used as a case study. Since dimensionless analysis is conducted, the results on natural frequencies obtained are valid for all thickness of the plate.

Table 2 illustrates the effects of the foundation parameter on natural frequency. In this case, consideration is given to (a) Elastic Winkler-type foundation ($g_s = 0, k_L = 10, 50, 100$), (b) Elastic Pasternak-type foundation ($k_L = 0, g_s = 10, 50, 100$) and (c) Two-parameter elastic foundation ($k_L = 50, g_s = 10, 50, 100$)

Parametric effect of nonlinear Winkler foundation is carried out as mentioned earlier using the finite difference method of numerical due to difficulty encountered when

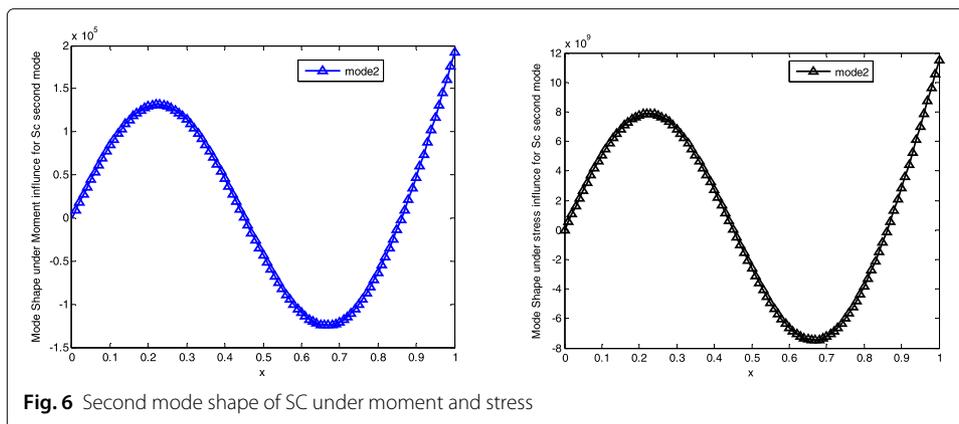


Table 2 Two-parameter foundation effect on first two natural frequencies of vibrating rectangular plate

Edge condition	Natural frequency	$k_L = 50, \lambda = 0.5$			$k_L = 50, \lambda = 1.5$		
		$g_s = 10$	$g_s = 50$	$g_s = 100$	$g_s = 10$	$g_s = 50$	$g_s = 100$
SS	Ω_1	17.3464	26.3758	34.4842	34.3159	39.653	45.4515
	Ω_2	46.9493	61.5091	75.8769	65.1907	76.3477	88.3341
SC	Ω_1	21.5364	30.1136	38.0430	37.2582	42.6778	48.5309
	Ω_2	56.5026	69.9870	83.7489	73.2330	84.0060	95.7134
SF	Ω_1	8.0257	7.0030	4.8888	24.7306	22.8481	19.1978
	Ω_2	22.8135	31.2448	39.1257	43.2468	48.6099	54.4463

analysing with ADM. Although, it is a universal behaviour of plate to be affected by characteristic of elastic foundation, comparing Table 2 to Table 1 indicates that properly chosen value of foundation stiffness allows for both foundation stiffness and plate stiffness effect to be comparable.

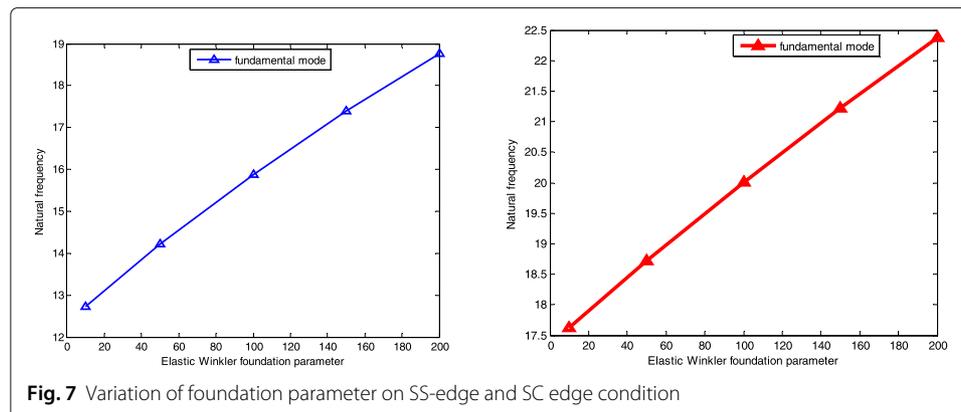
It is clearly shown in all cases that increasing the foundation stiffness results into higher value of natural frequencies. Moreover, it is also observed that the effect of the difference in natural frequencies is more significant for higher mode of the vibrating rectangular plate. Figure 7 shows the graphical relationship of increasing foundation stiffness on the natural frequency. The graph depicts the linear relationship between stiffness increment and natural frequency.

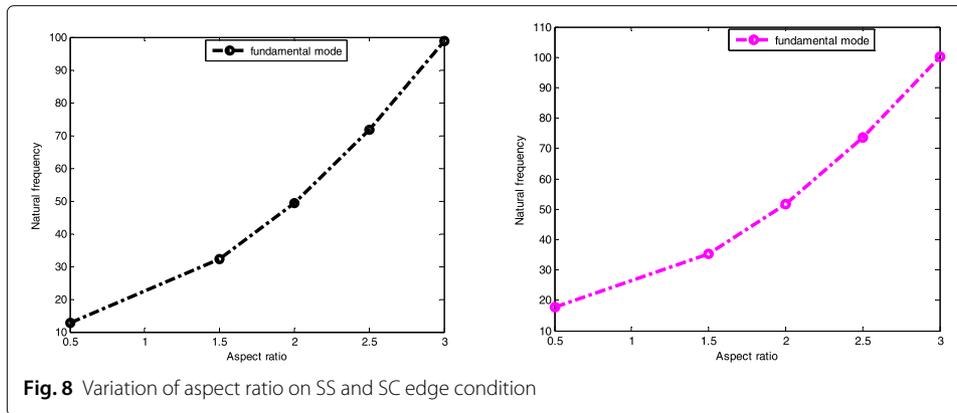
Variation of aspect ratio on natural frequency

Considering the effect of change in the aspect ratio, recall that uniform thickness thin plate is the model. It is clearly seen that the reduction in the value of the aspect ratio (negative value) yields the same result as the positive variation of the aspect ratio increase in aspect ratio either in all cases results into increase in natural frequency of plate (Fig. 8). This phenomenon is a result of uniform thickness assumption in the model. The variation reduction in thickness has no effect.

Submerging of plate in fluid

Determining the effect of water in contact with plate involves incorporation of fluid pressure parameters into the model as stated earlier in the assumptions and using





the parameters stated for the simulation. In order to illustrate the validity of the proposed analytical solutions, the numerical computation of the isotropic rectangular plate submerged in water is done and compared with those available from previous studies as shown in Table 3. The following parameters are adopted for the simulation according to [9].

$A = 0.20165\text{ m}$, $B = 0.65\text{ m}$, $h = 0.963\text{ mm}$. Reservoir dimension $1.3\text{ m} \times 0.55\text{ m} \times 0.8\text{ m}$.

Water has a noticeable effect on the characteristic behaviour of plate in contact with fluid. Table 4 depicts the natural frequency of plate when submerged and immersed in fluid. The fluid under consideration here is water; it is globally known that natural frequency becomes lower when the vibrating plate is in contact with fluid. The behaviour is referred to as added virtual mass effect. The behaviour is the result of the transfer of vibration by plate in contact with water in motion thereby increasing the kinetic energy of the total system. It is found that the presence of water has most significant effect on the fundamental mode regardless of the boundary condition of the rectangular plate.

Conclusion

In this study, the investigation of dynamic response of isotropic rectangular plates resting on Winkler and Pasternak foundations is presented. The coupled governing nonlinear partial differential equation is transformed to nonlinear ordinary differential equation using the Galerkin method of separation. The coupled nonlinear ordinary differential equations have been solved using the Adomian decomposition method. Moments and stress are determined. The accuracies of the obtained analytical solutions were ascertained with the numerical and analytical methods reported in the literature, and the results are confirmed to be in good agreement. The obtained analytical solutions were used to examine the effects of foundation parameter, aspect ratio and presence of fluid in contact with plate. From the parametric studies, the following observations were established. Increases in the elastic foundation parameter increases the natural frequency. Increases in aspect ratio increases the natural frequency. Submerging the plate in fluid

Table 3 Validation of solution for fluid-plate interaction

Natural frequencies of simply supported plate submerged in fluid			
Mode	Haddara and Cao [35]	Present	Percentage error
1	28.72	22.118	22
2	117.725	108.993	7
3	154.51	156.596	1.35

Table 4 Comparative study of the effect of plate resting on fluid to other medium

Edge condition/ natural frequency	Simply supported (SS)			Simply Supported clamped (SC)			Simply supported free (SF)		
	Fluid with Winkler and Pasternak foundations	Submerged in fluid	Immersed in fluid	Fluid with Winkler and Pasternak foundations	Submerged in fluid	Immersed in fluid	Fluid with Winkler and Pasternak foundations	Submerged in fluid	Immersed in fluid
Ω_1	21.7888	0.0000	18.0570	25.6351	0.0000	22.2613	17.4634	14.1776	8.5408
Ω_2	53.0700	43.1951	48.6995	62.0728	53.5723	58.1018	53.8591	57.0730	26.5863
Ω_3	103.0190	95.7538	98.3596	117.2008	110.6347	112.8977	100.4398	113.8003	61.3445

lowers the natural frequency. ‘The major benefit of ADM is the ease of application to various types of differential and integral equations, linear or nonlinear, homogeneous or inhomogeneous, with constant coefficients or with variable coefficients; likewise, it is capable of greatly reducing the size of computational work while still maintaining high accuracy of the numerical solution.’ [36].

Abbreviations

$\frac{d}{dx}$: First-order differential operator with respect to x ; ν : Poisson’s ratio of isotropic plate; ρ : Mass density; ω : Dimensionless natural frequency; D : Flexural rigidity of isotropic plate; h : Plate thickness; a, b : Dimension of the plate; w : Transverse deflection; x, y : Rectangular space Cartesian coordinate along the length of thin plate

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Authors’ contributions

SA analysed and interpreted the data required for this research; he also put up the manuscript. AY proofread, checked and edited the entire work. MG wrote the manuscript. All authors read and approved the final manuscript.

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Availability of data and materials

The datasets used and/or analysed during the current study are available from the corresponding author on reasonable request.

Competing interests

The authors declare that they have no competing interests.

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